

Camp elèctric (I)

Introducció

Càrrega elèctrica: Magnitud física fonamental (longitud, temps, massa,...)

De dos signes

Conservada (energia)

Quantitzada $\pm e$ (sempre múltiples $\pm Ne$)

Partícules fonamentals de la matèria

Electró $-e$

$$m_e = 9.1096 \times 10^{-31} \text{ kg}$$

Protó e

$$\begin{aligned} m_p &= 1.673 \times 10^{-27} \text{ kg} \\ &= 1836 m_e \end{aligned}$$

Unitat de càrrega: el coulomb (C)

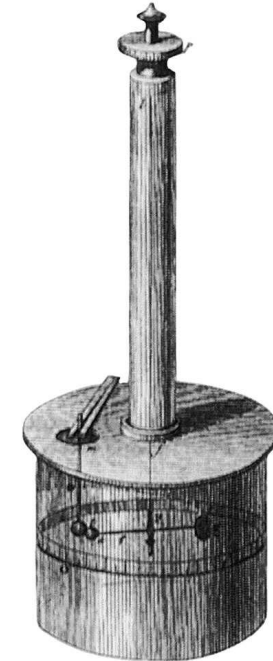
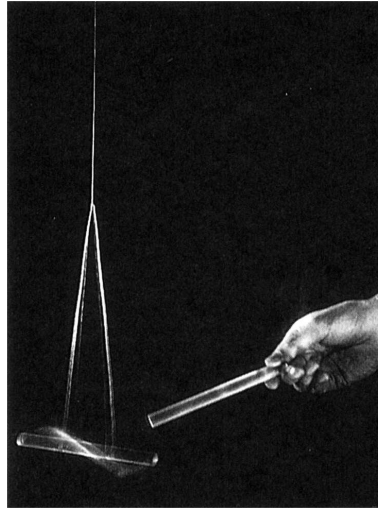
$$e = 1.6022 \times 10^{-19} \text{ C}$$

$$1 \text{ C} = 6.2414 \times 10^{18} e$$

Exemple: a una moneda de coure de 3 g hi ha $\sim 8.24 \times 10^{23}$ electrons, i la seva càrrega total és $-1.32 \times 10^5 \text{ C}$

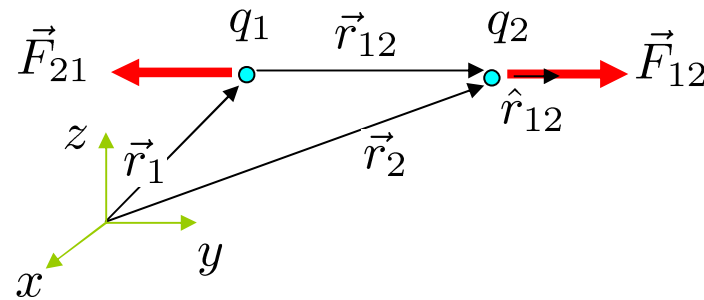
La llei de Coulomb

Els cossos amb càrregues de signe contrari s'atreuen i els que tenen càrregues del mateix signe es repel·leixen



The force exerted by one point charge on another acts along the line between the charges. It varies inversely as the square of the distance separating the charges and is proportional to the product of the charges. The force is repulsive if the charges have the same sign and attractive if the charges have opposite signs.

Càrregues 'puntuals'



Vector de 1 a 2

$$\vec{r}_{12} = \vec{r}_2 - \vec{r}_1$$

$$\hat{r}_{12} = \frac{\vec{r}_{12}}{r_{12}} \quad \text{vector unitari}$$

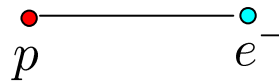
$$\vec{F}_{12} = k \frac{q_1 q_2}{r_{12}^2} \hat{r}_{12}$$

$$\vec{F}_{21} = k \frac{q_1 q_2}{r_{21}^2} \hat{r}_{21}$$

$$9 \cdot 10^9 \frac{\text{N}}{\text{m}^2 \text{C}^2}$$

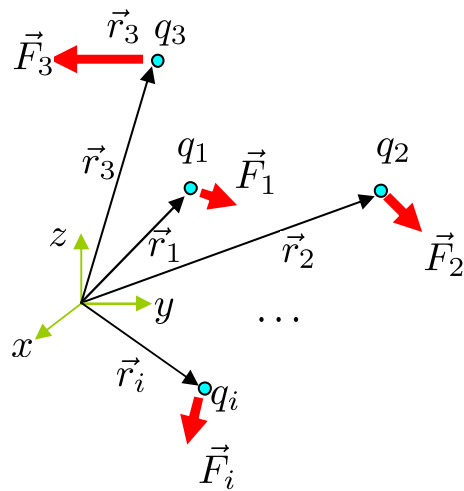
Davant la força de Coulomb la gravetat es menyspreable

$$\vec{F}_{12}^{(g)} = G \frac{m_1 m_2}{r_{12}^2} \hat{r}_{12}$$



$$\frac{F_{12}}{F_{12}^{(g)}} = \frac{ke^2}{Gm_p m_e} = 2.27 \times 10^{39}$$

Més càrregues (superposició)

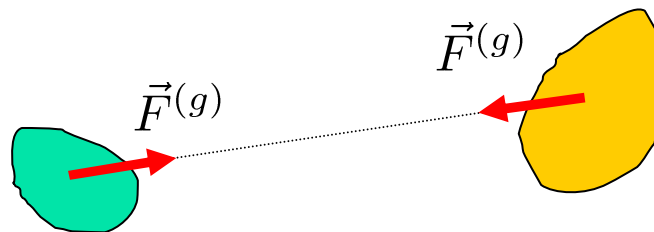


$$\vec{F}_i = \vec{F}_{1i} + \vec{F}_{2i} + \vec{F}_{3i} + \dots$$

$$= \sum_{\substack{j=1 \\ (j \neq i)}}^N k \frac{q_j q_i}{r_{ji}^2} \hat{r}_{ji}$$

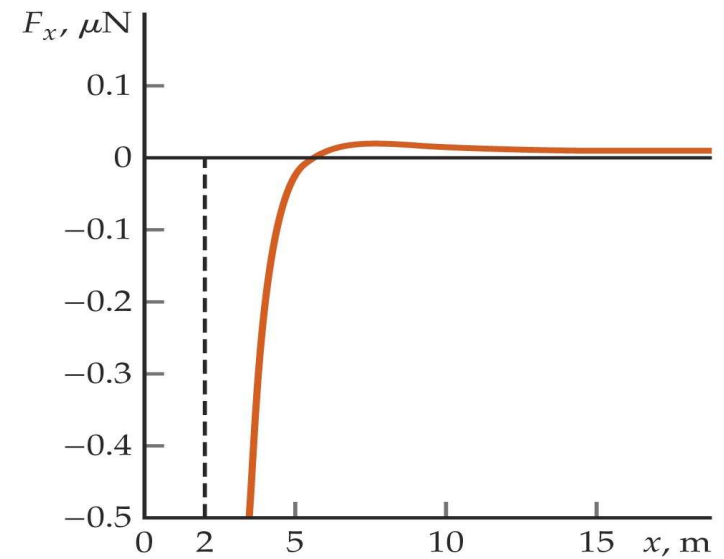
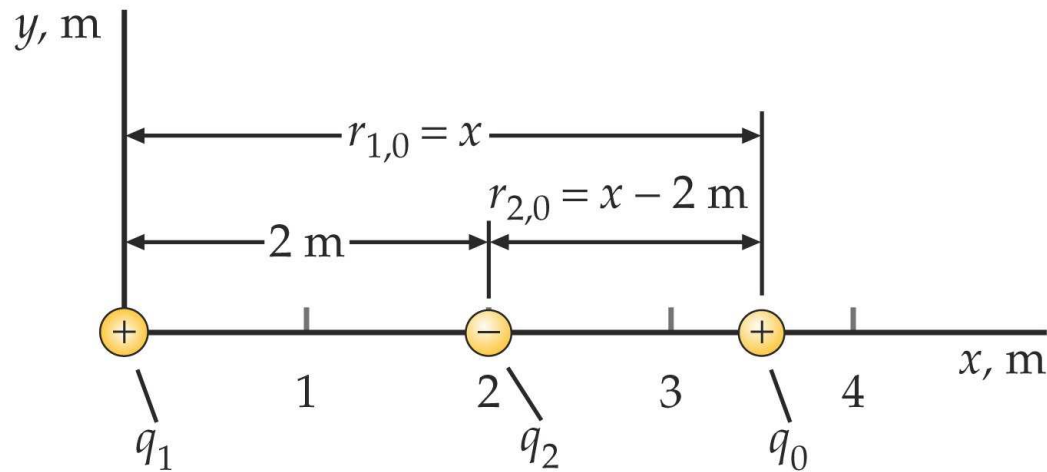
Compensació de càrregues al mon macroscòpic

Només forces gravitatòries





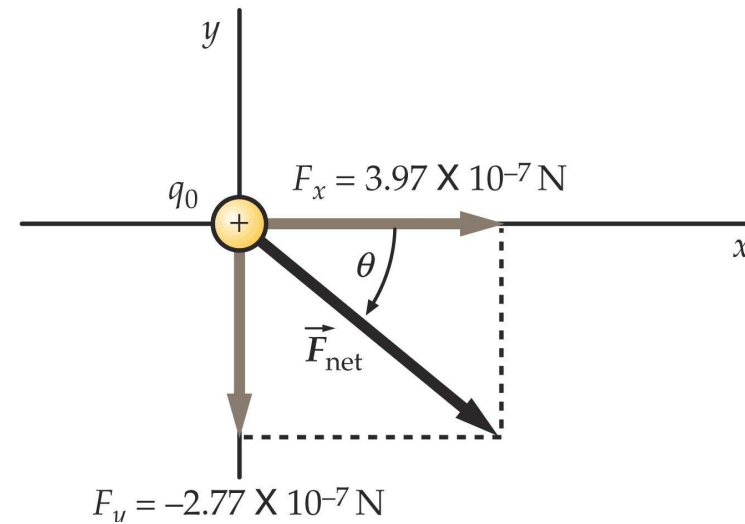
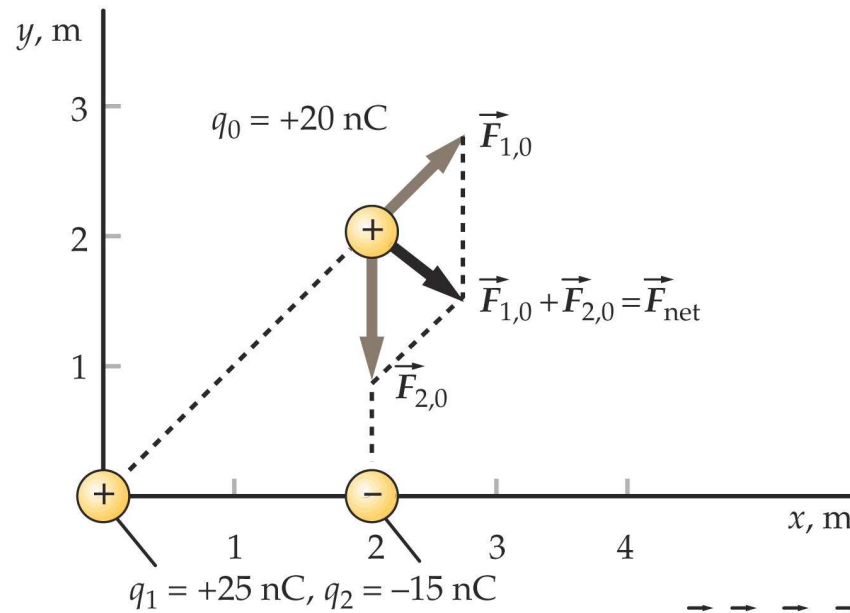
Exemple de força neta en una dimensió



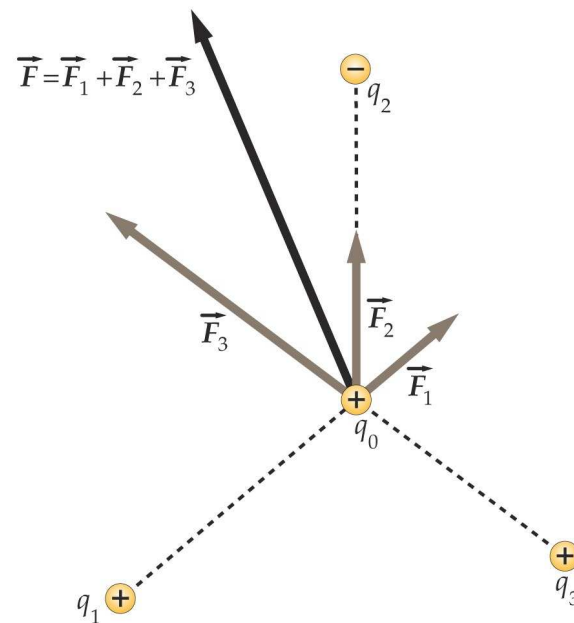
$$q_1 = 25 \text{ nC}$$

$$q_2 = -10 \text{ nC}$$

Exemple de força neta en dues dimensions



I en tres dimensions

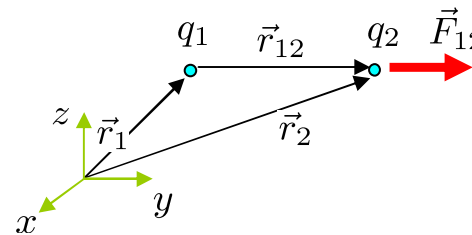
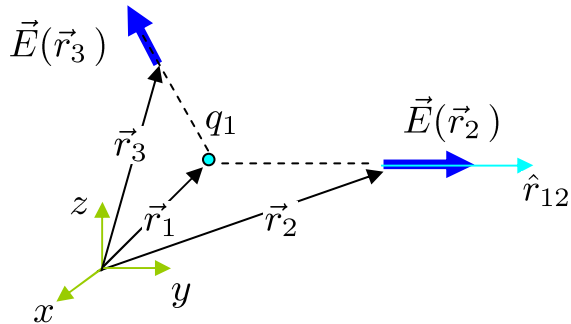


El camp elèctric

Interpretació de la llei de Coulomb:

q_1 crea un *camp vectorial* a tot l'espai $\vec{E}(\vec{r})$

q_2 experimenta una força $\vec{F}_{12}$ degut al camp de q_1



$$\vec{F}_{12} = q_2 \vec{E}(\vec{r}_2)$$

$$\vec{E}(\vec{r}_2) = k \frac{q_1}{r_{12}^2} \hat{r}_{12}$$

$$\vec{E}(\vec{r}_2) = k \frac{q_1}{r_{12}^3} \vec{r}_{12}$$

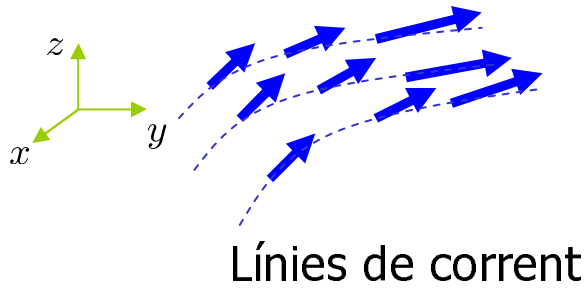
vector unitari de $\vec{r}_1$ a $\vec{r}_2$

$$\hat{r}_{12} = \frac{\vec{r}_{12}}{r_{12}}$$

Càrrega que crea el camp (font) i càrrega de prova
Es pot reinterpretar canviant els papers

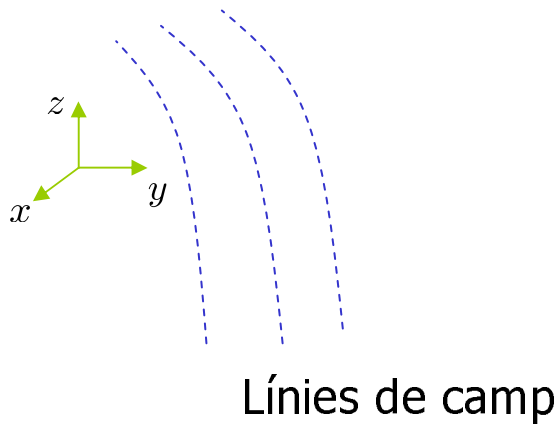


Camps vectorials: un vector a cada punt de l'espai



Camp de velocitats $\vec{v}(\vec{r})$

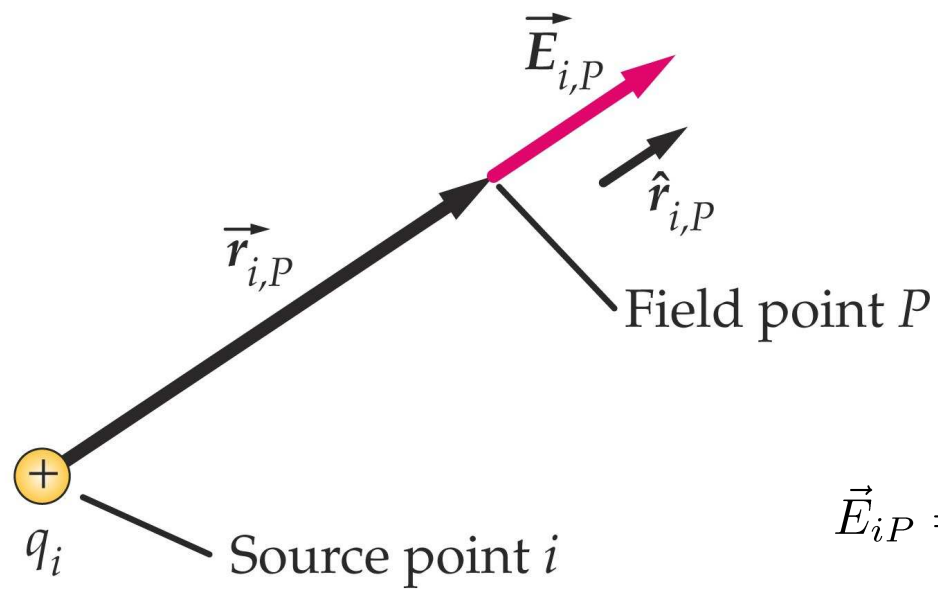
Camps escalars: un escalar a cada punt de l'espai



Camp de temperatures $T(\vec{r})$

La permitivitat ϵ_0

$$k = \frac{1}{4\pi\epsilon_0}$$



$$\vec{E}_{iP} = \frac{q_i}{4\pi\epsilon_0 r_{iP}^2} \hat{r}_{iP} = \frac{q_i}{4\pi\epsilon_0 r_{iP}^3} \vec{r}_{iP}$$

Exemple: Si una càrrega de 5 nC experimenta una força de $2 \times 10^{-4} \text{ N } \vec{u}_x$, quin es el camp elèctric a seva posició ?

$$\vec{E} = \frac{\vec{F}}{q_0} = \frac{2 \times 10^{-4} \text{ N } \vec{u}_x}{5 \times 10^{-9} \text{ C}} = 4 \times 10^4 \frac{\text{N}}{\text{C}} \vec{u}_x$$

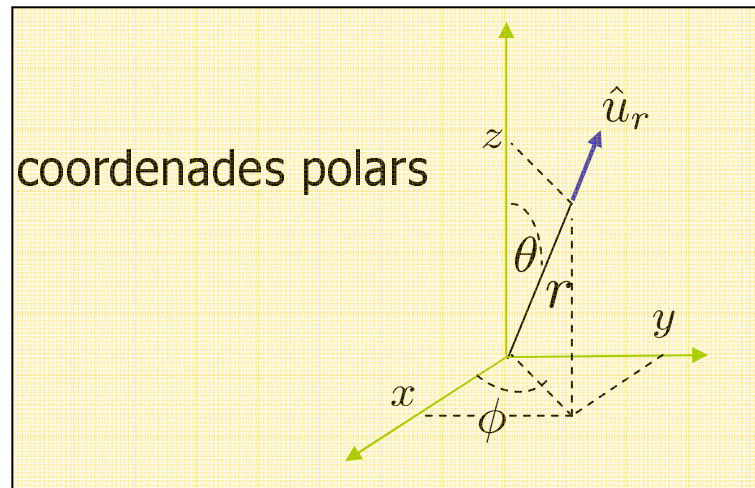
↙ Unitat de camp elèctric

TABLE 21-2

Some Electric Fields in Nature

	$E, \text{N/C}$
In household wires	10^{-2}
In radio waves	10^{-1}
In the atmosphere	10^2
In sunlight	10^3
Under a thundercloud	10^4
In a lightning bolt	10^4
In an X-ray tube	10^6
At the electron in a hydrogen atom	6×10^{11}
At the surface of a uranium nucleus	2×10^{21}

Ordres de magnitud



Camp d'una càrrega puntual

Diagram showing the electric field vector $\vec{E}(\vec{r})$ originating from a point charge Q at the origin of a 3D coordinate system. The vector points radially outwards at a distance r . The axes are labeled x , y , and z .

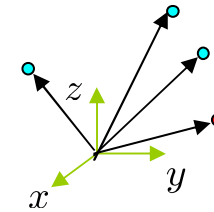
$$\vec{E}(\vec{r}) = \frac{Q}{4\pi\epsilon_0 r^2} \hat{u}_r$$



Més d'una càrrega (superposició)

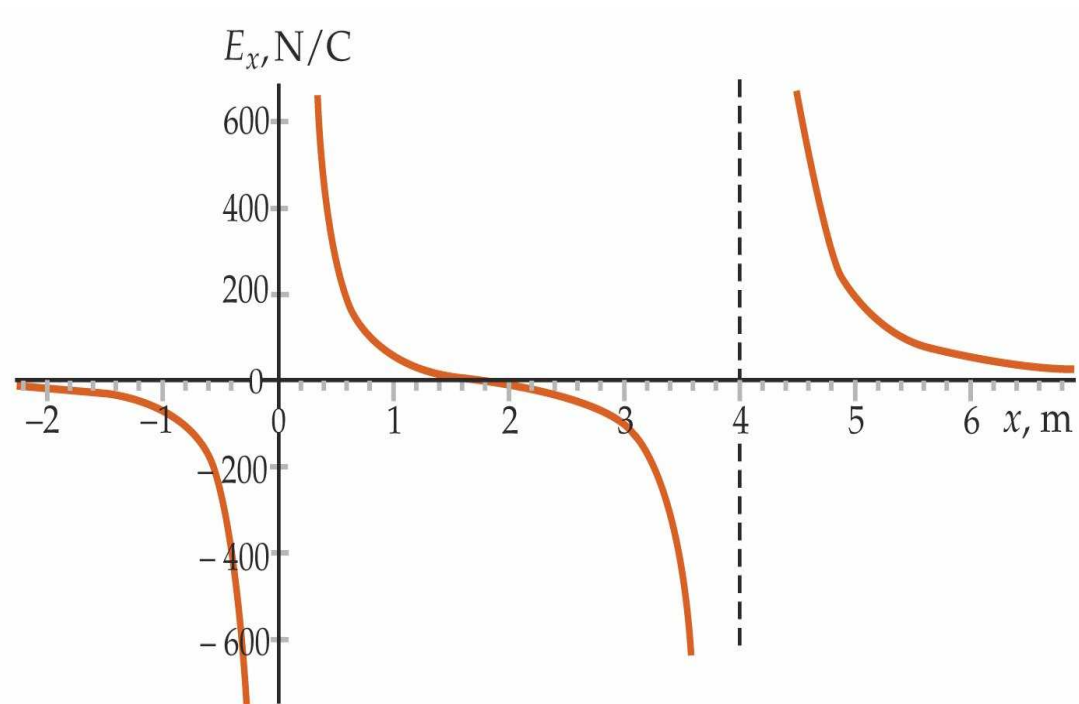
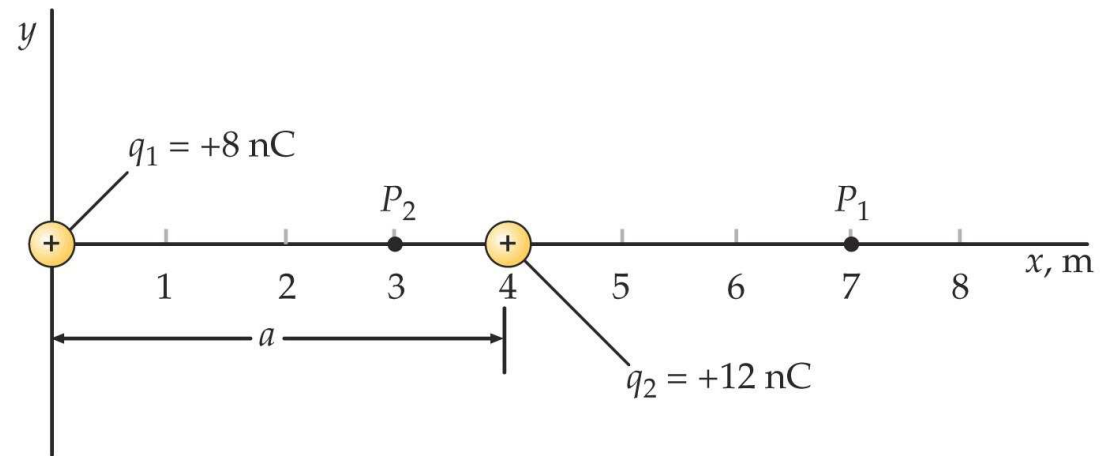
Càrregues font $(q_1, \vec{r}_1), (q_2, \vec{r}_2), \dots (q_N, \vec{r}_N)$

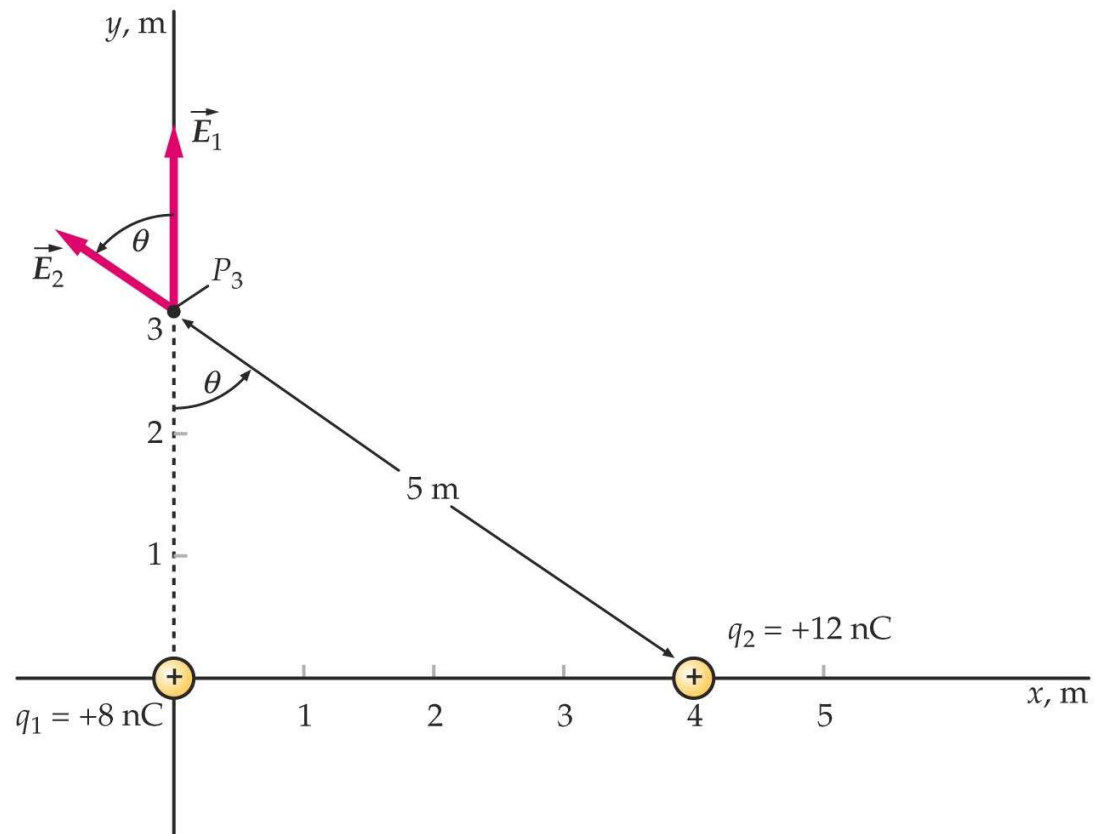
Càrrega prova $(q_p, \vec{r}_p)$

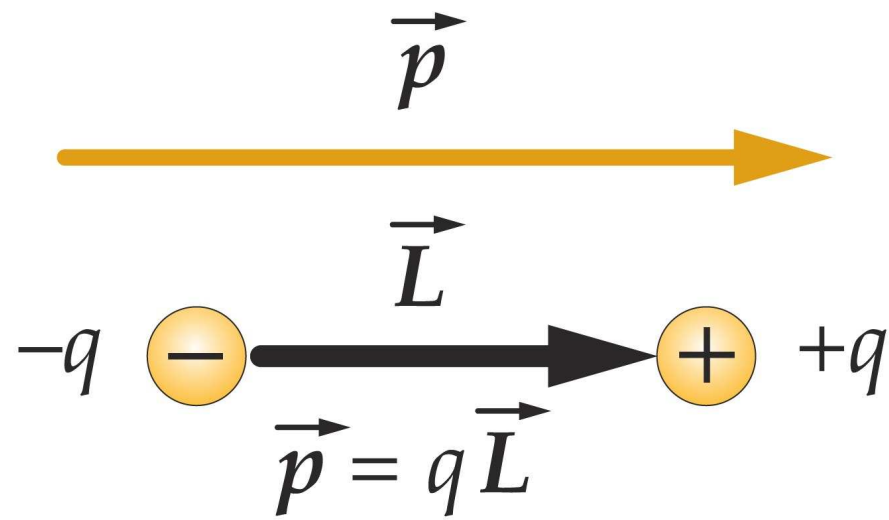


$$\vec{E}_p = \sum_{i=1}^N k \frac{q_i}{r_{ip}^2} \hat{r}_{ip} = \sum_{i=1}^N k \frac{q_i}{r_{ip}^3} \vec{r}_{ip}$$

$$\vec{F}_p = q_p \vec{E}_p$$



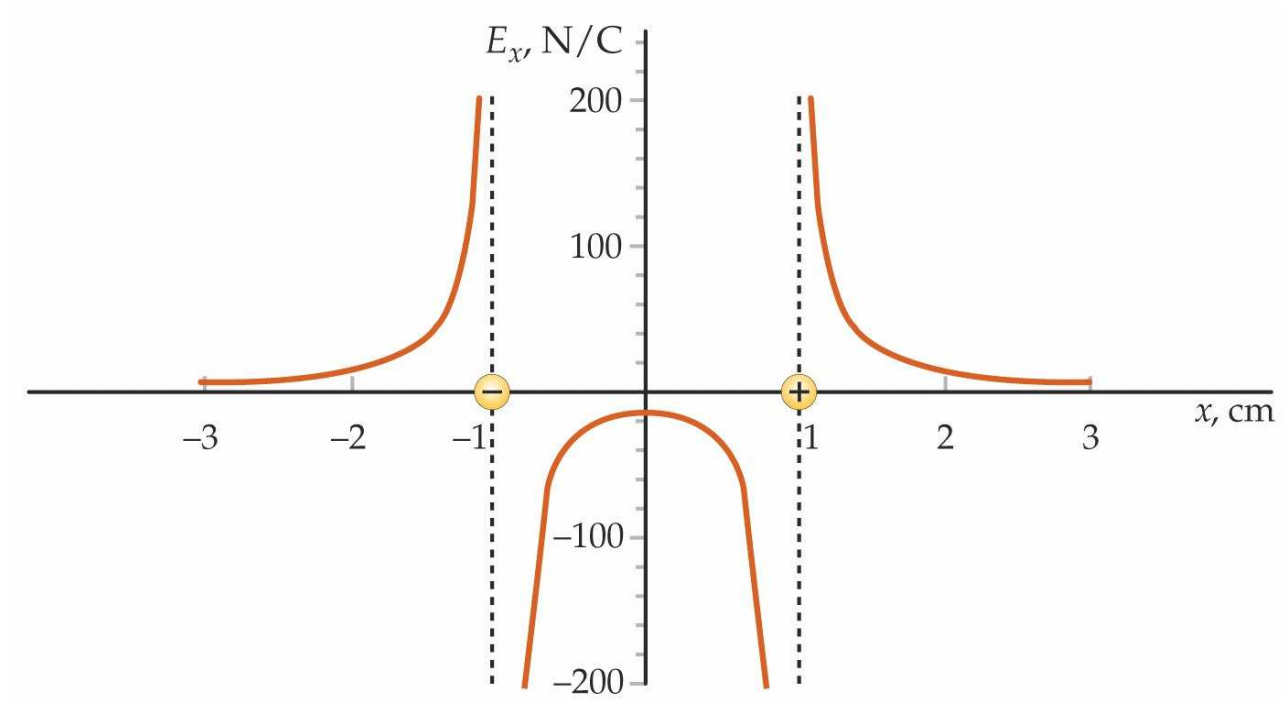




$$\vec{p} = q\vec{L}$$

21-9

DEFINITION—ELECTRIC DIPOLE MOMENT



A llargues distàncies
 $x \rightarrow \infty$

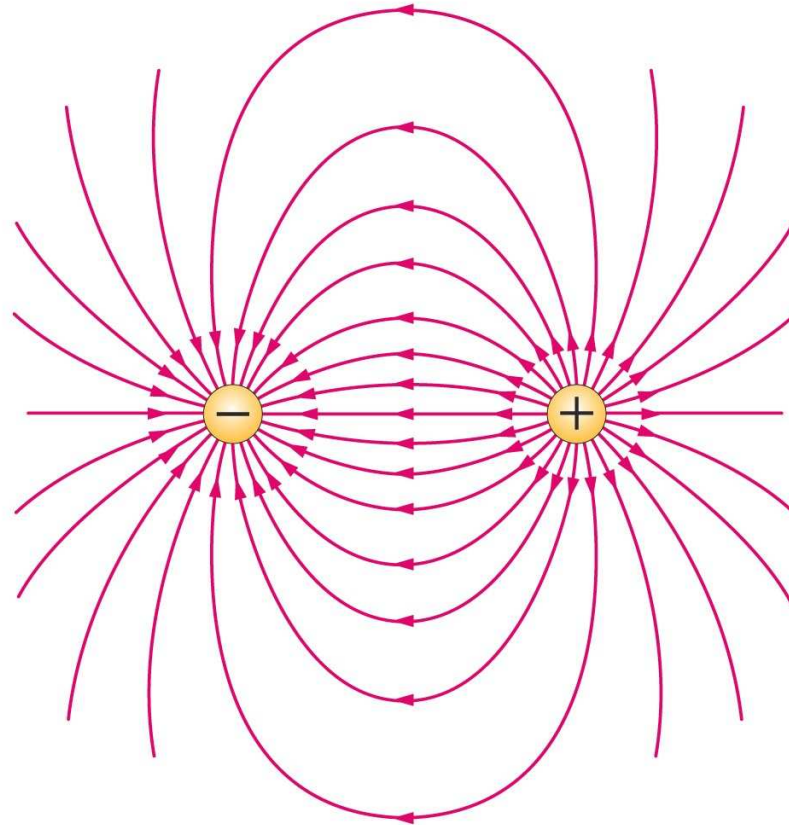
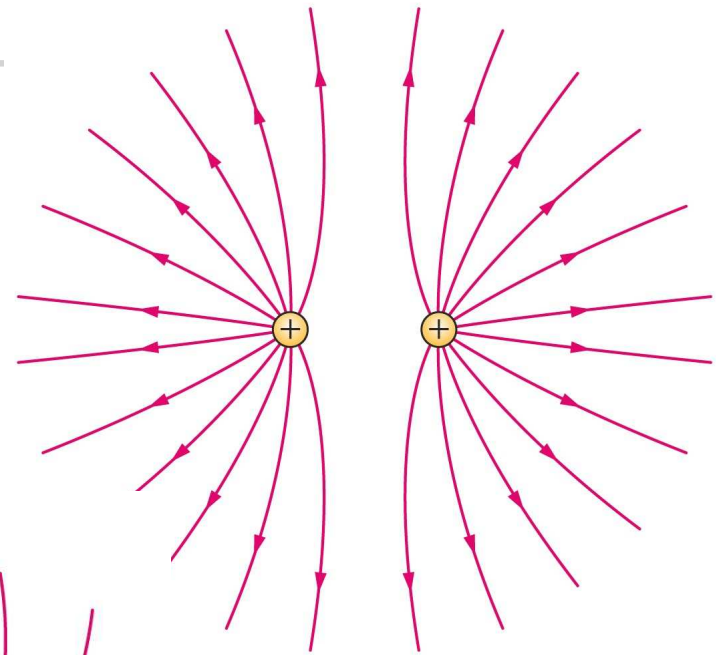
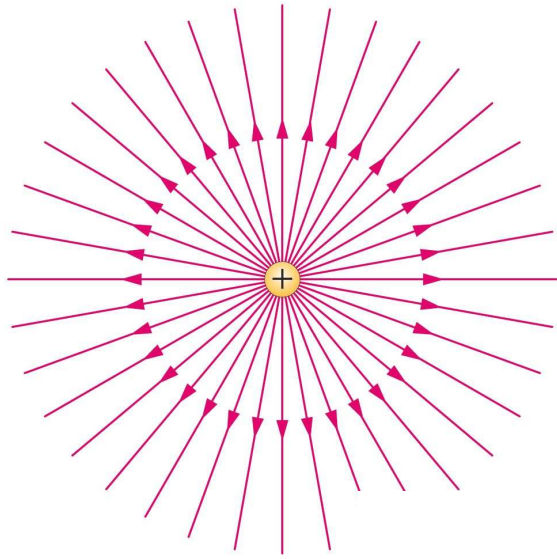
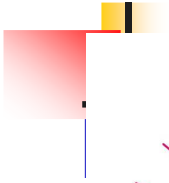
$$E_x = \frac{2kp}{|x|^3}$$

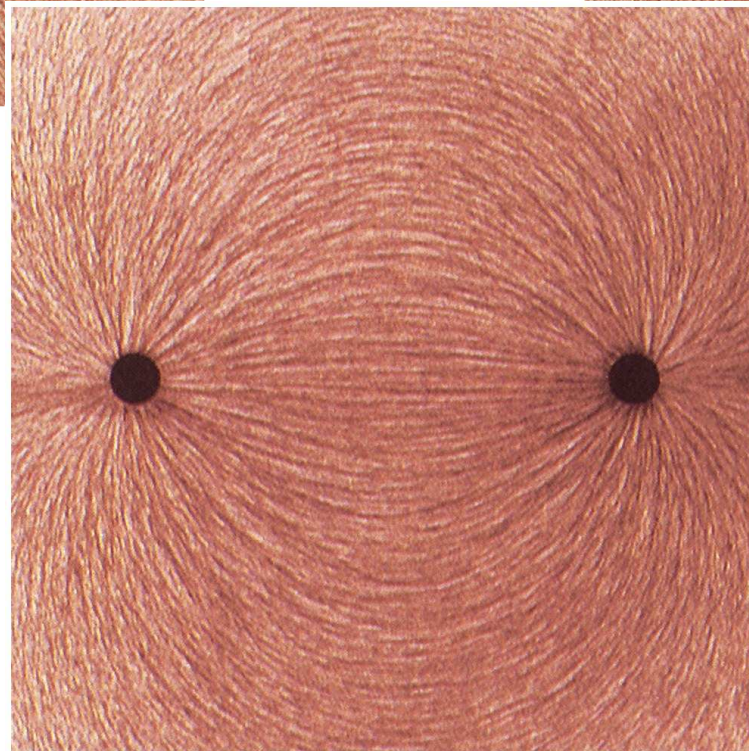
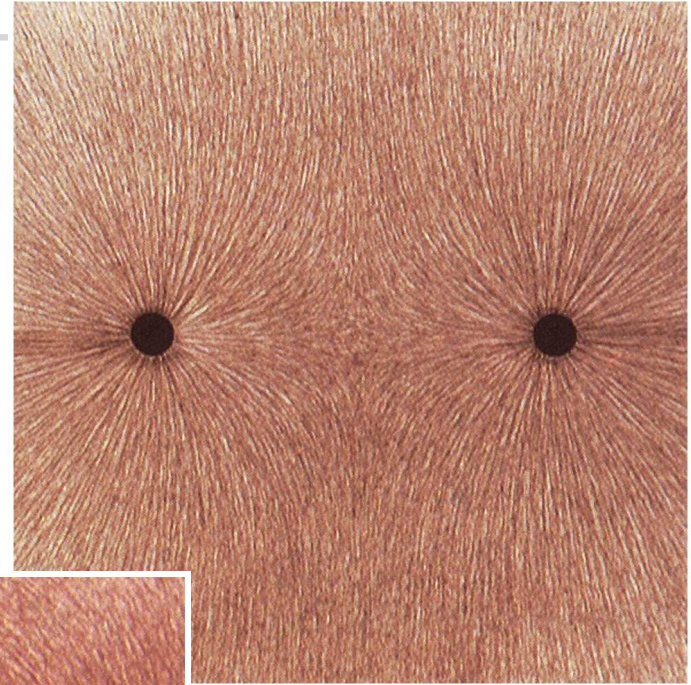
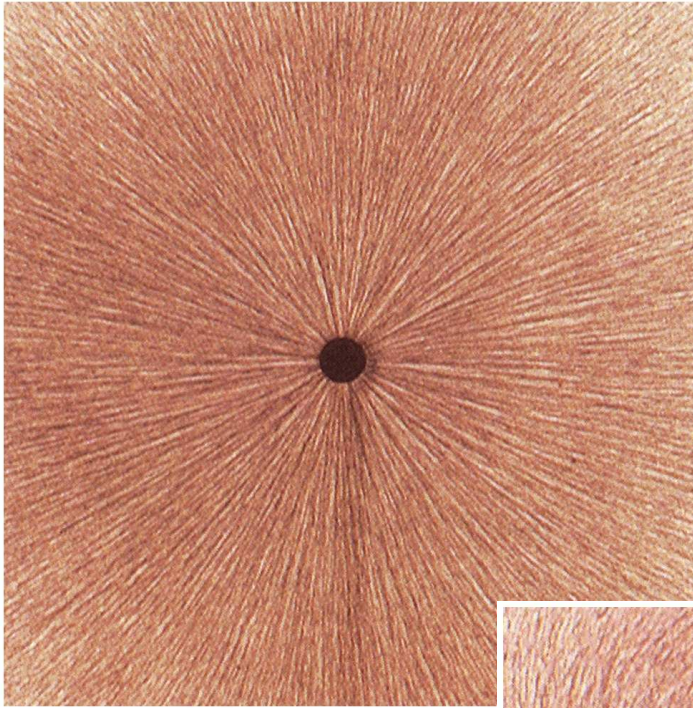


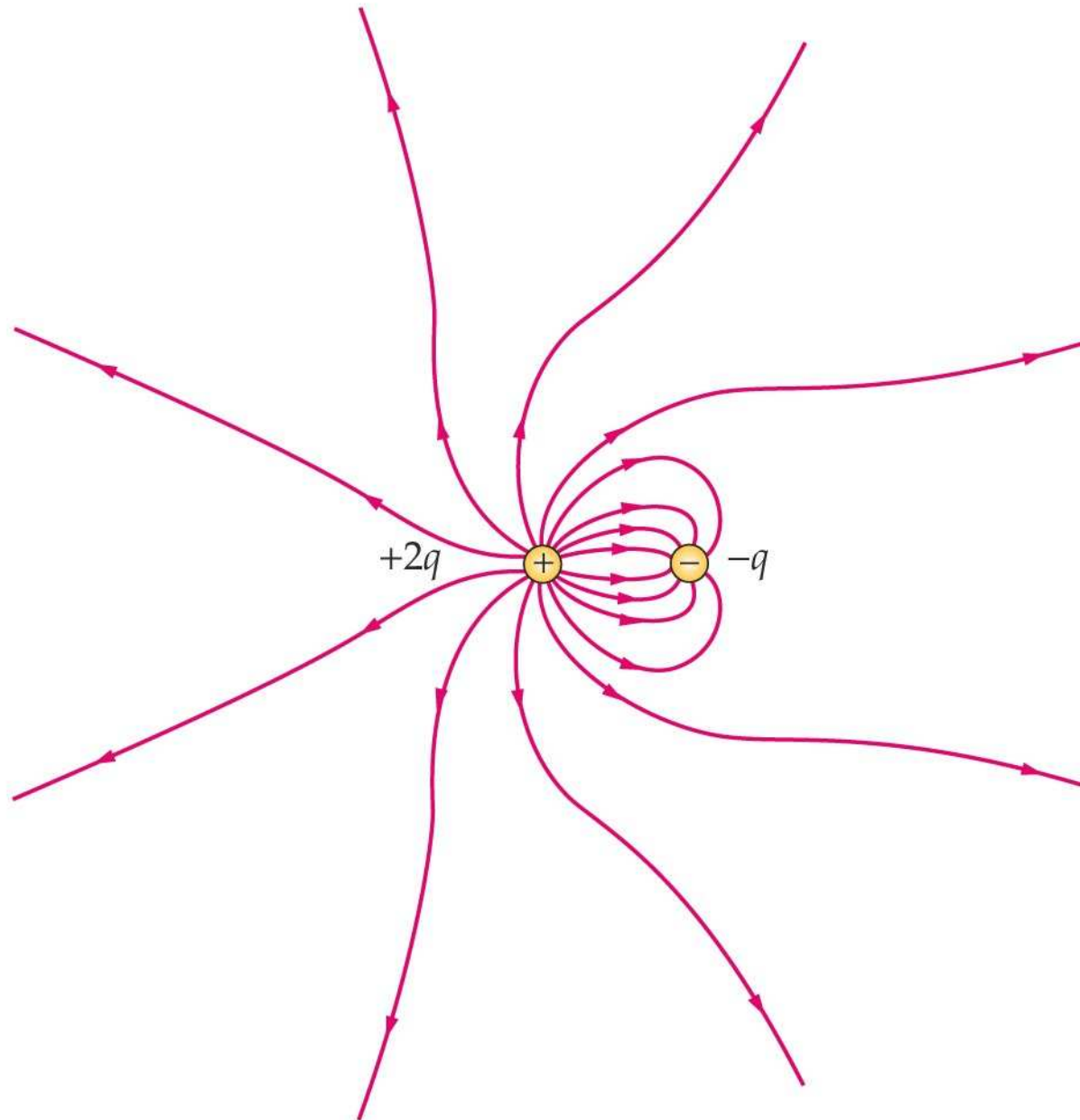
Línies de camp (elèctric)

1. Electric field lines begin on positive charges (or at infinity) and end on negative charges (or at infinity).
2. The lines are drawn uniformly spaced entering or leaving an isolated point charge.
3. The number of lines leaving a positive charge or entering a negative charge is proportional to the magnitude of the charge.
4. The density of the lines (the number of lines per unit area perpendicular to the lines) at any point is proportional to the magnitude of the field at that point.
5. At large distances from a system of charges with a net charge, the field lines are equally spaced and radial, as if they came from a single point charge equal to the net charge of the system.
6. Field lines do not cross. (If two field lines crossed, that would indicate two directions for $\vec{E}$ at the point of intersection.)

RULES FOR DRAWING ELECTRIC FIELD LINES

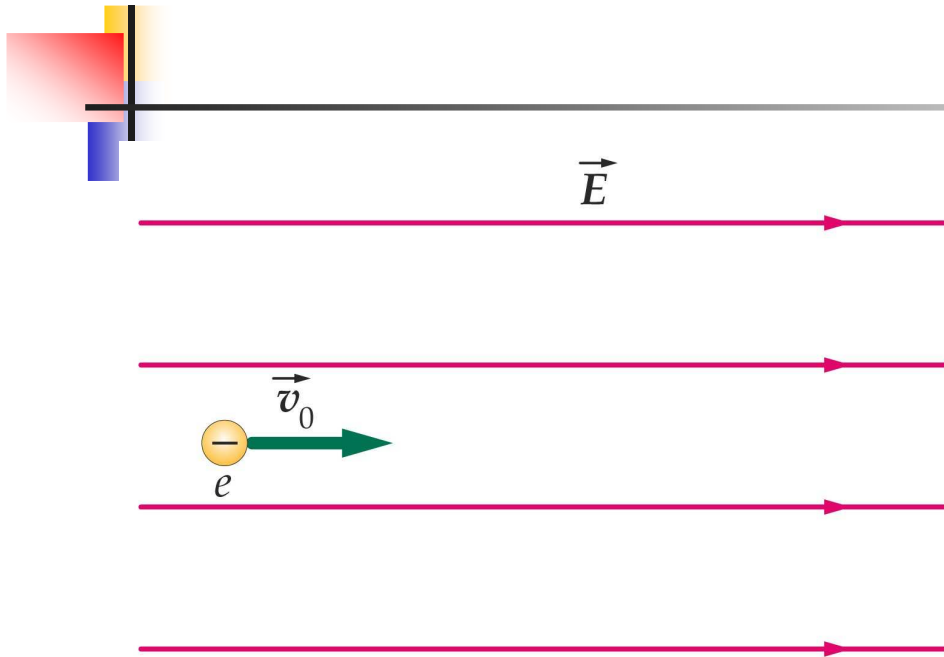






Moviments de càrregues dins camps electric





$$a_x = -\frac{eE}{m}$$

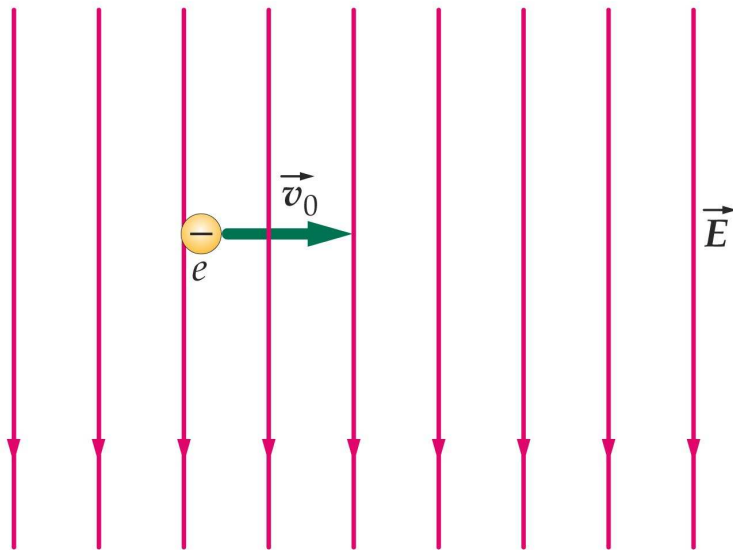
$$a_y = 0$$

$$v_x(t) = v_0 + a_x t$$

$$v_y(t) = 0$$

$$x(t) = x_0 + v_0 t + \frac{1}{2} a_x t^2$$

$$y(t) = y_0$$



$$a_x = 0$$

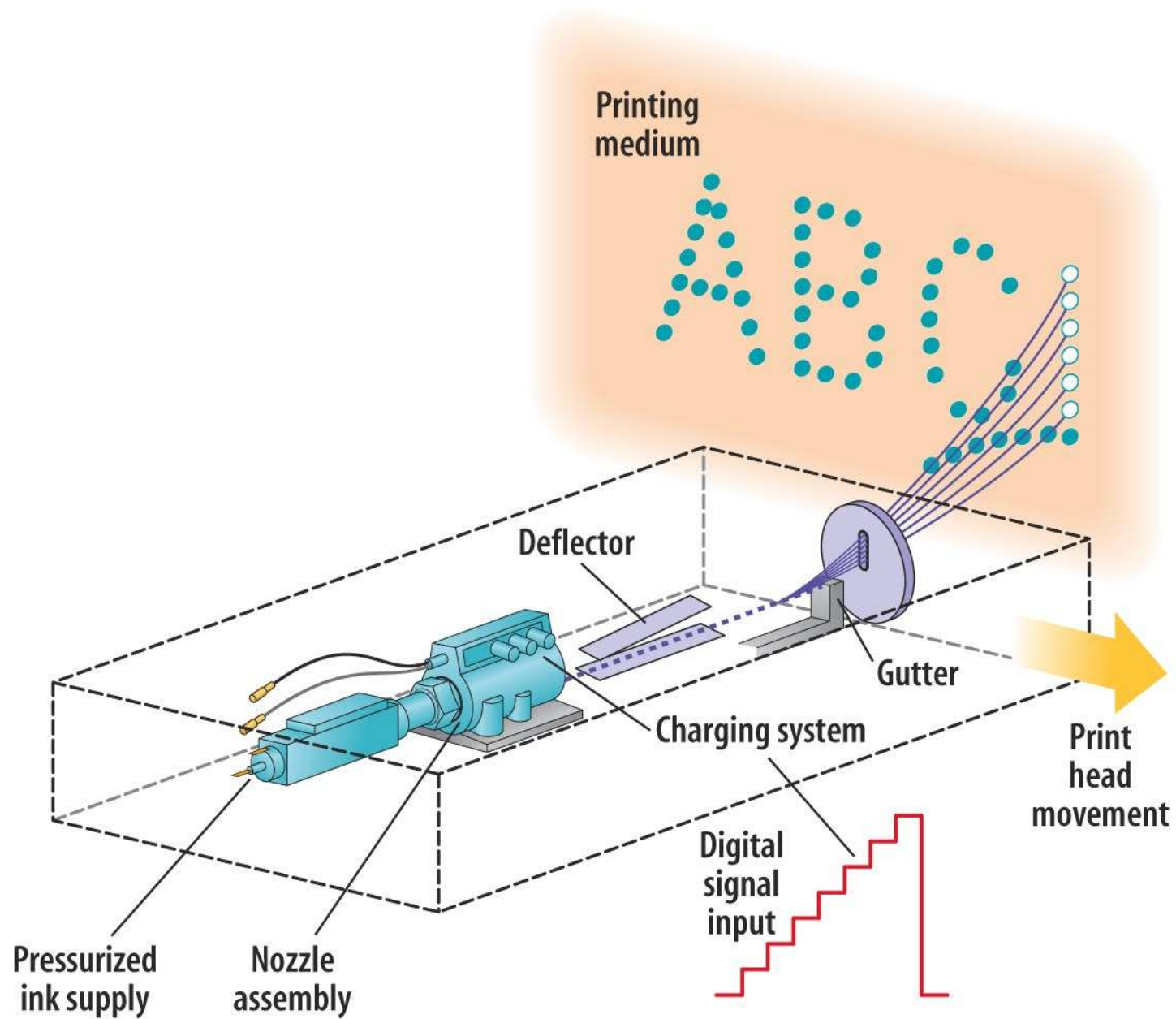
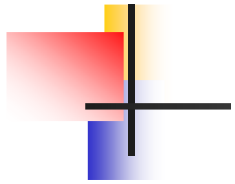
$$a_y = -\frac{eE}{m}$$

$$v_x(t) = v_0$$

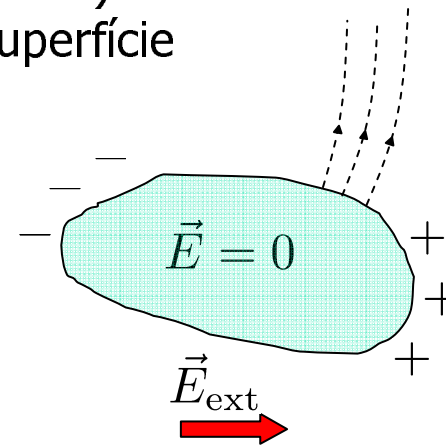
$$v_y(t) = a_y t$$

$$x(t) = x_0 + v_0 t$$

$$y(t) = y_0 + \frac{1}{2} a_y t^2$$

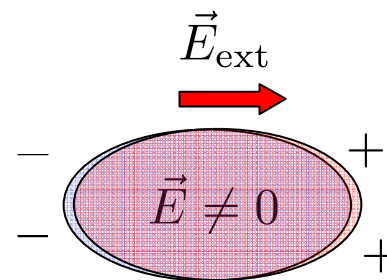
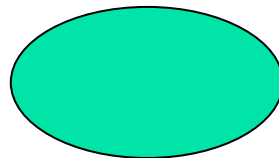


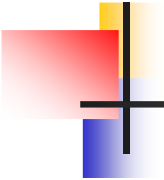
Conductors: moviment lliure de càrregues a l'interior
càrregues a la superfície (equilibri)
línies de camp normals a la superfície



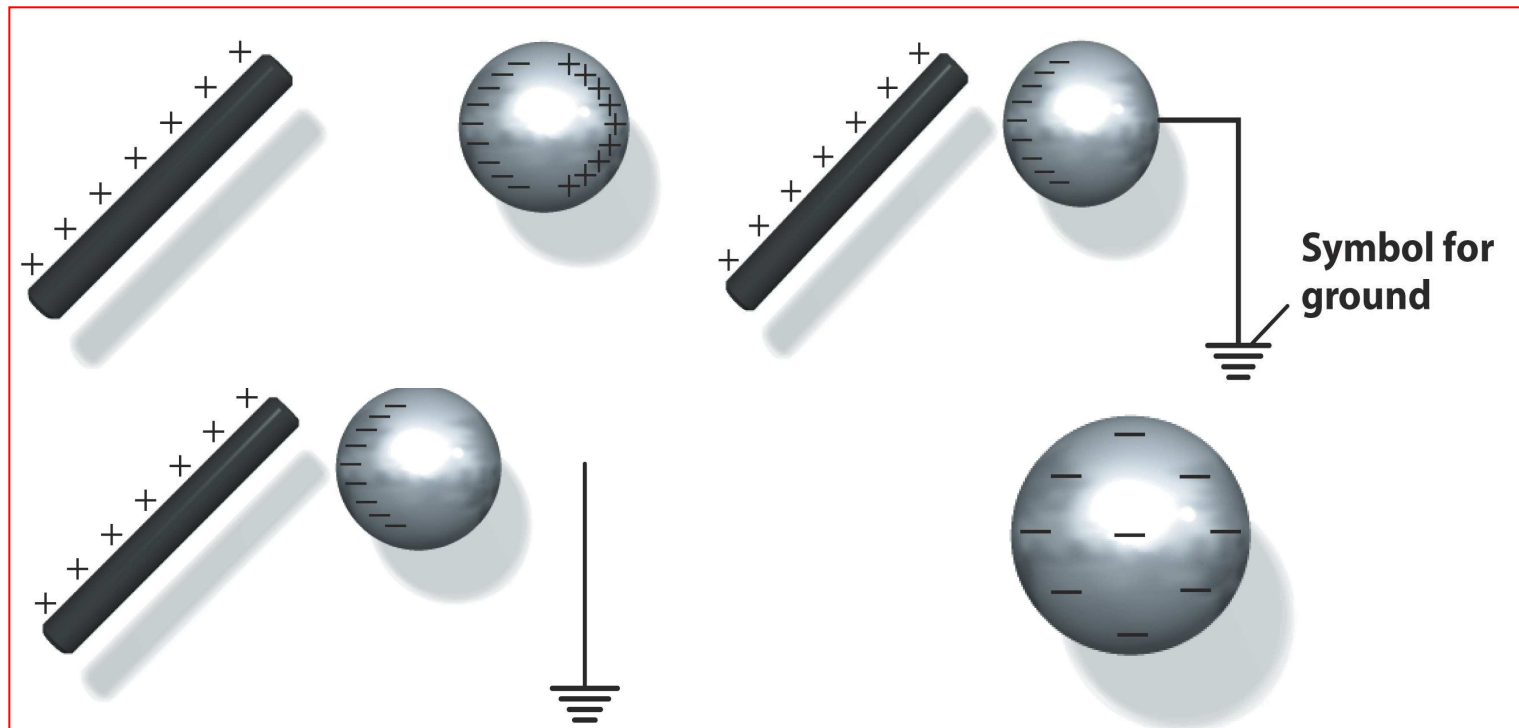
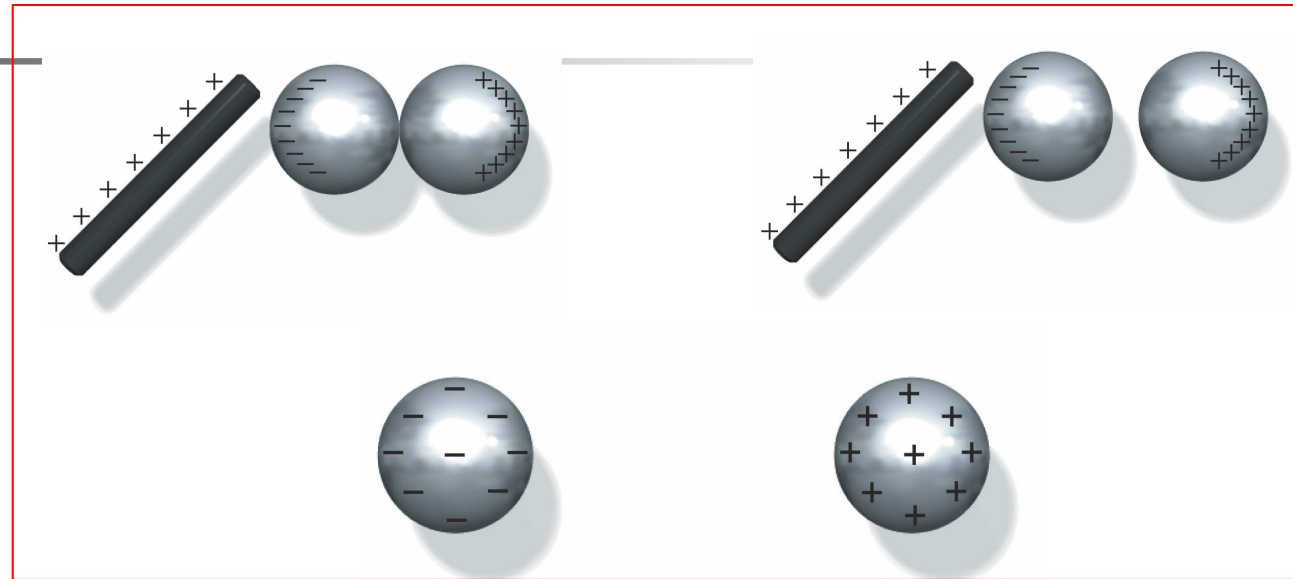
Aïllants: (dielèctrics)

Moviments limitats

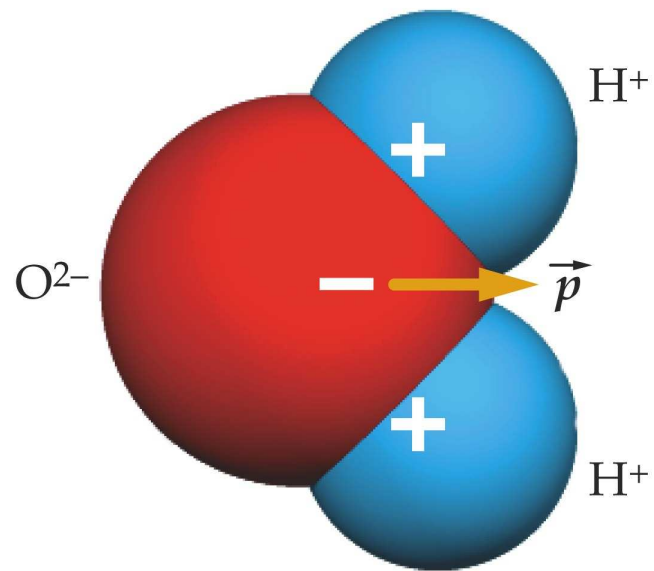


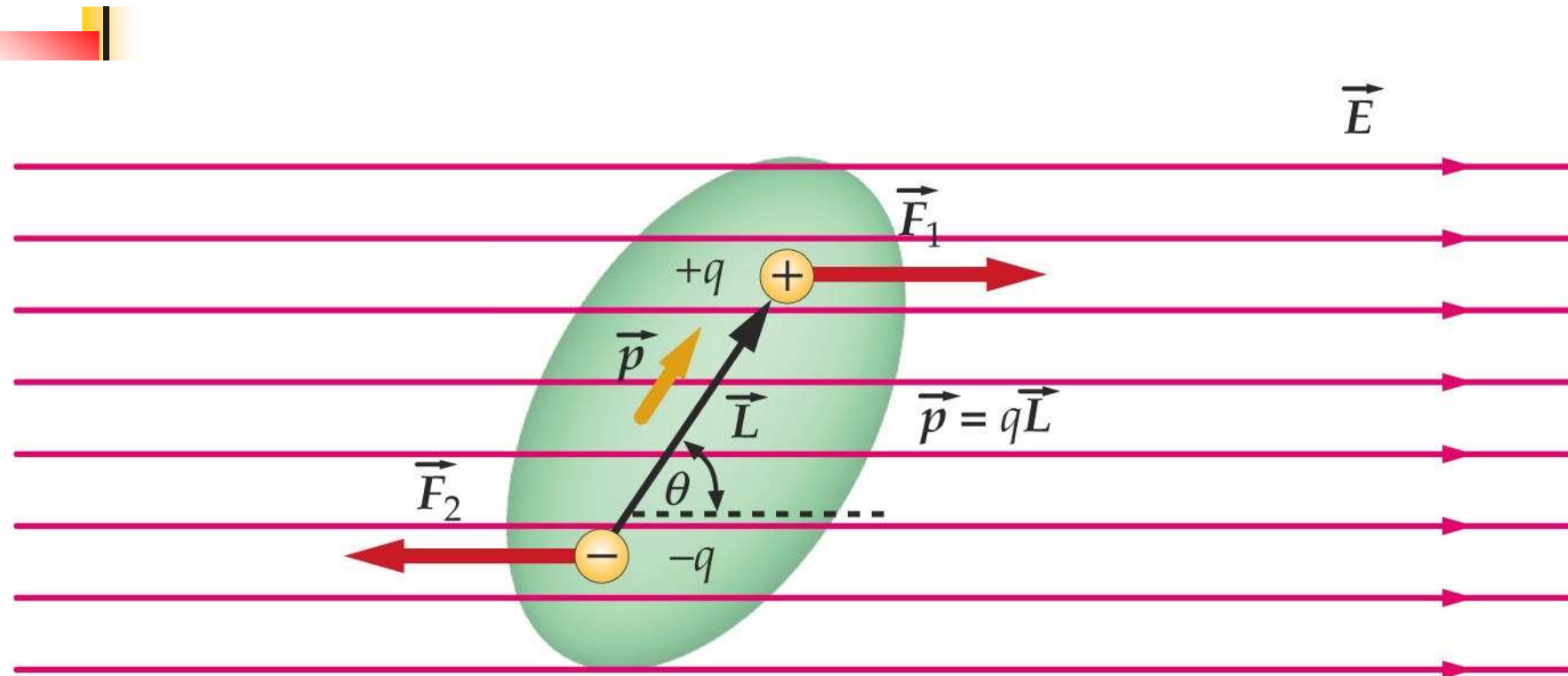


Càrrega de conductors per inducció



Electric Dipoles in Electric Fields



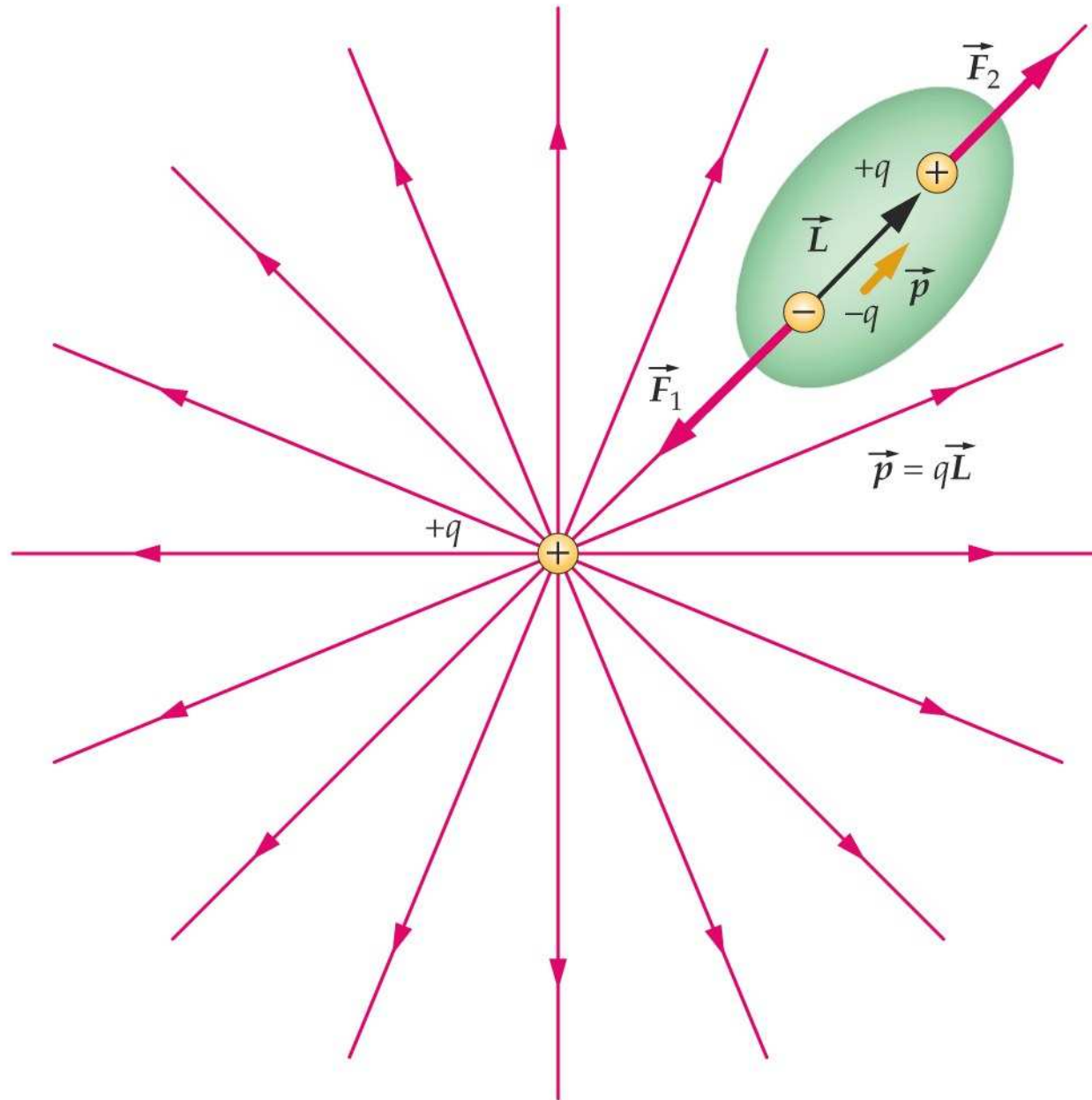


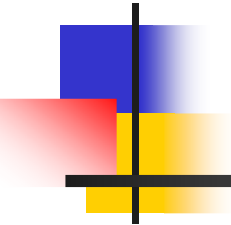
$$\tau = \vec{p} \times \vec{E}$$

$$U = -pE \cos \theta = -\vec{p} \cdot \vec{E}$$

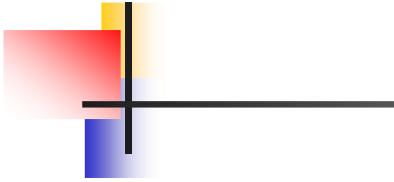
21-12

POTENTIAL ENERGY OF A DIPOLE IN AN ELECTRIC FIELD

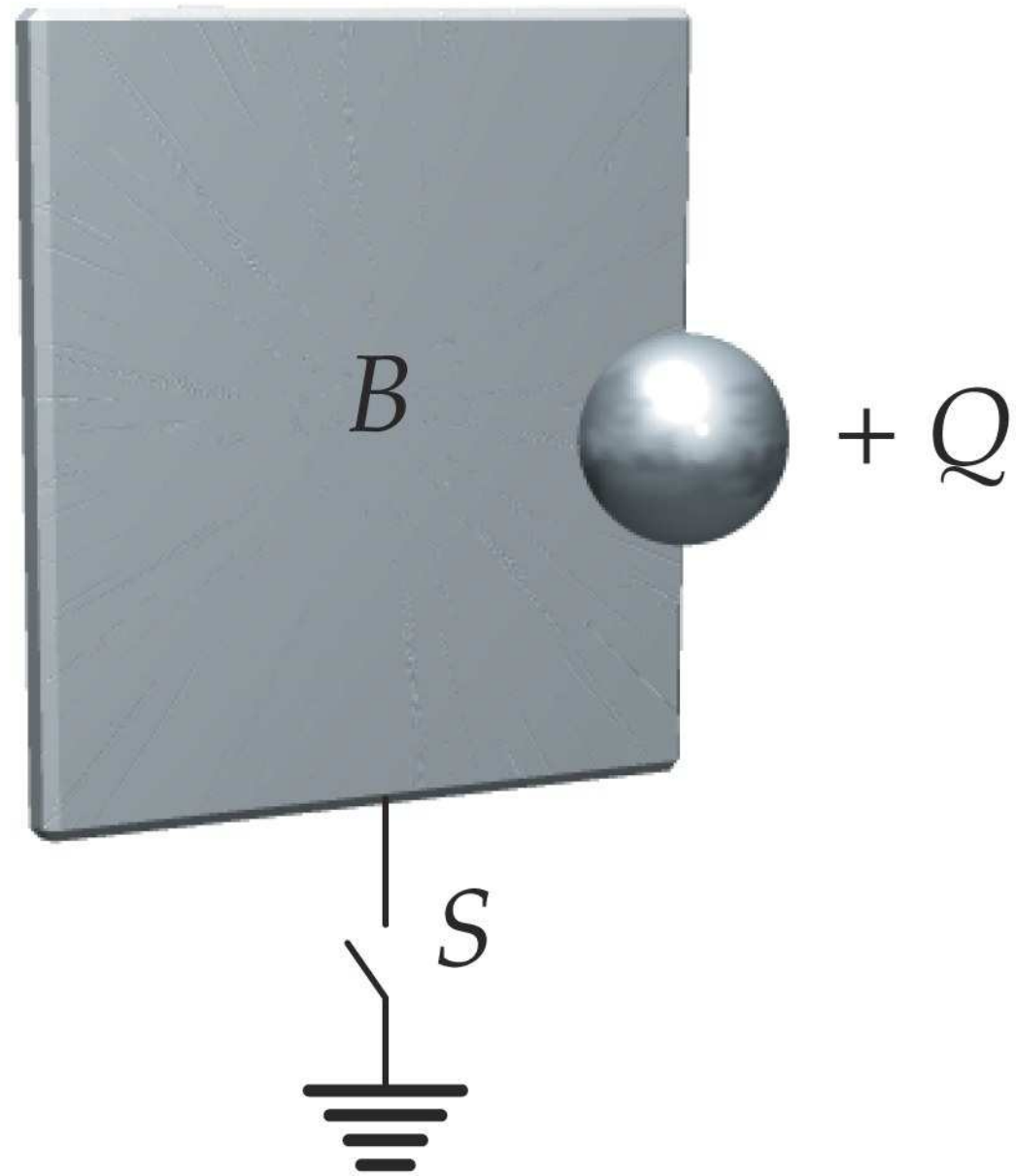


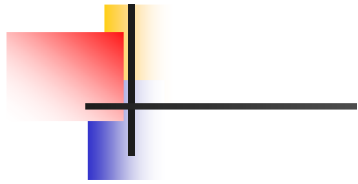


Problemes



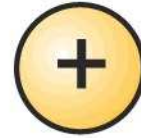
21.3



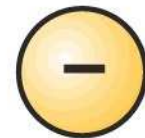


21.6

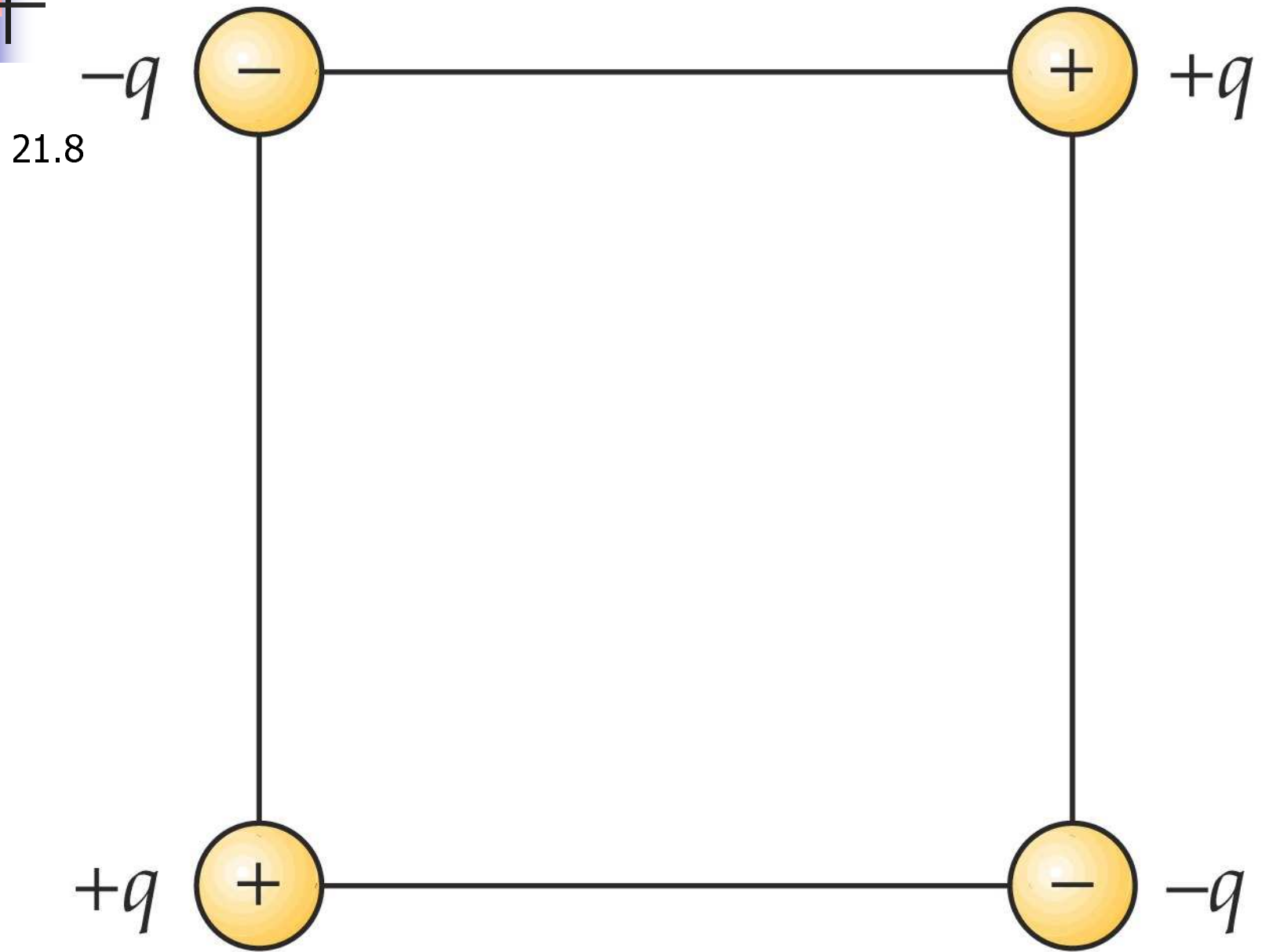
$+q$



$+Q$

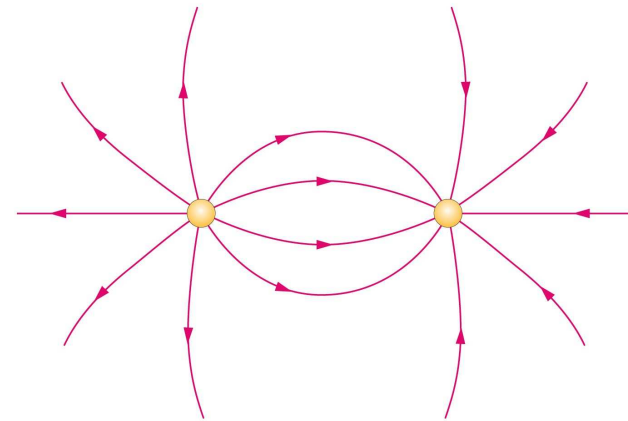
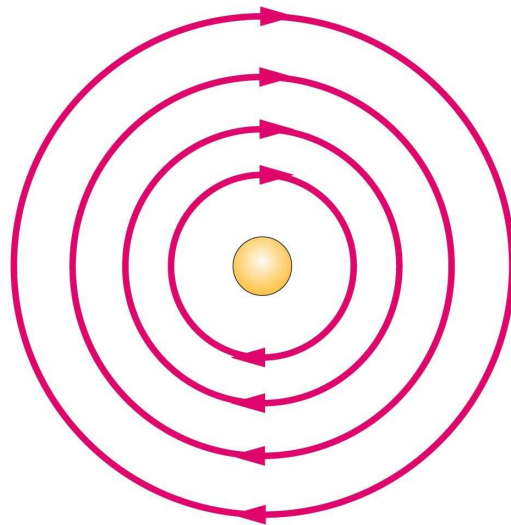
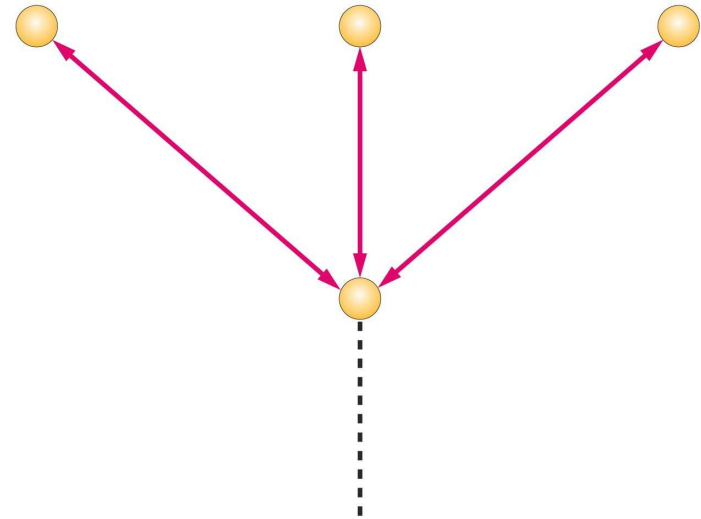
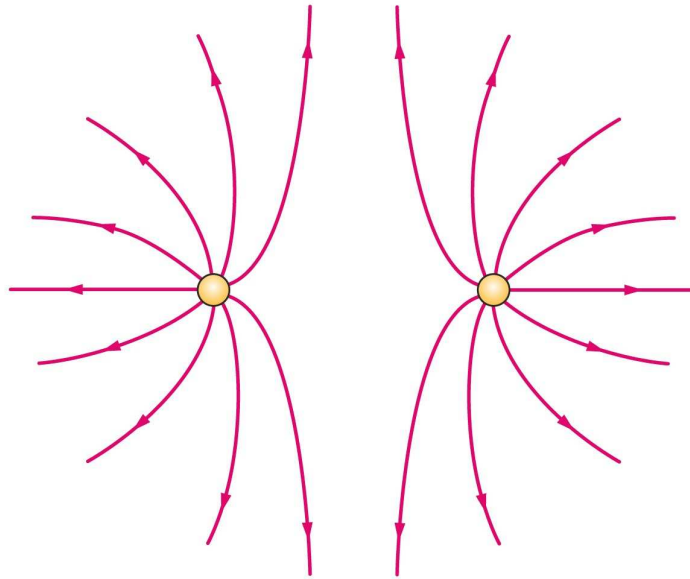


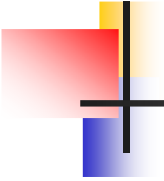
$-Q$



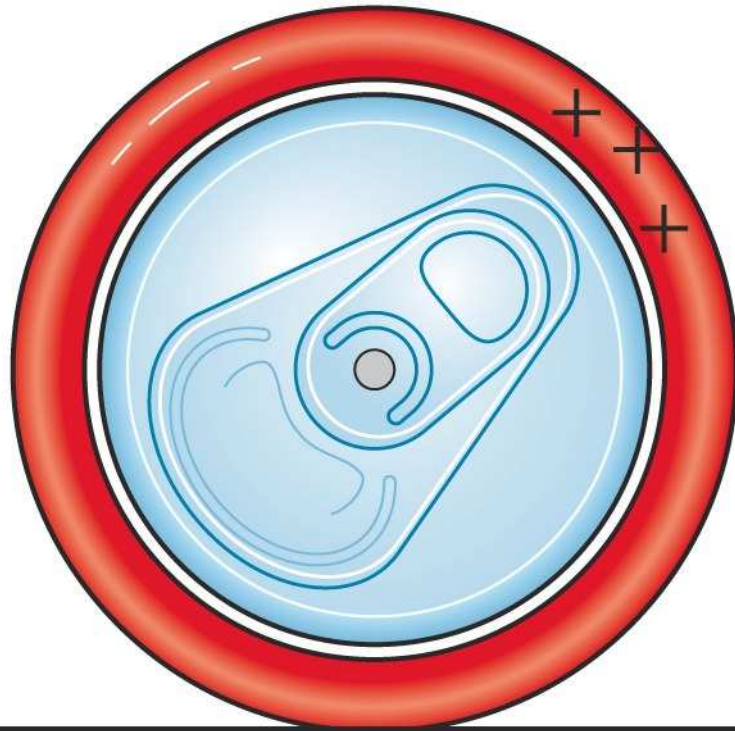


21.14





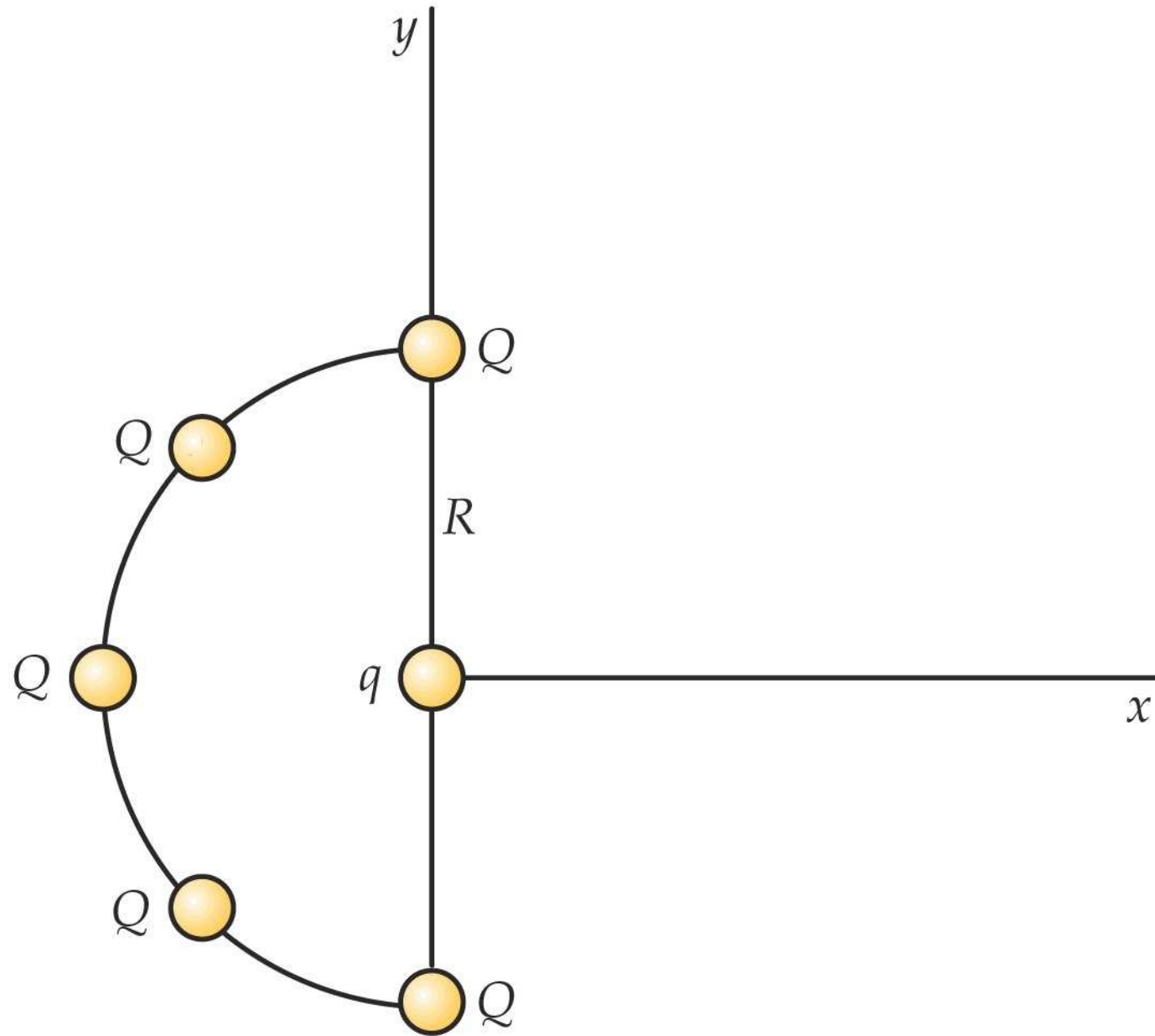
21.23



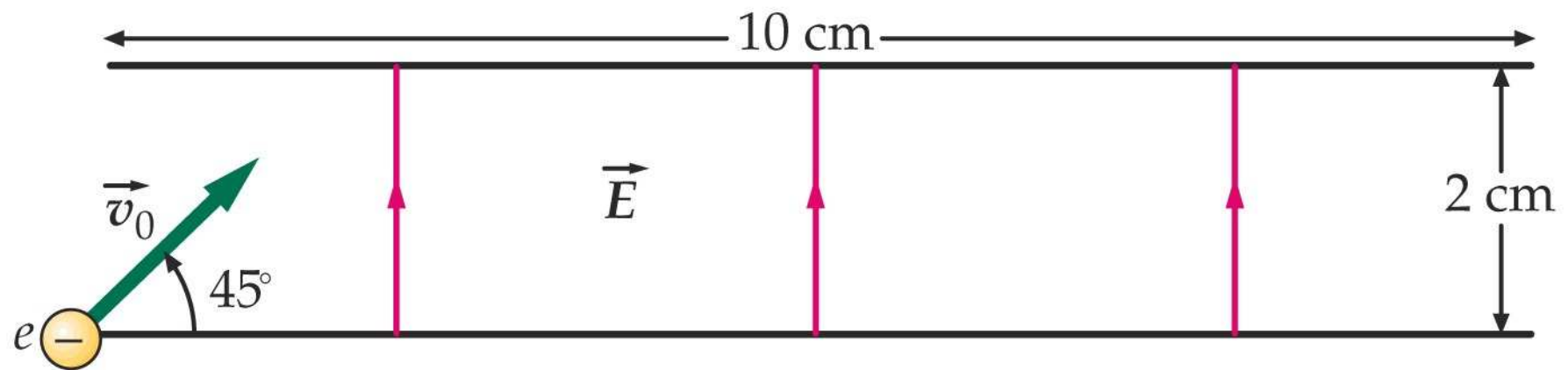
Soda
can



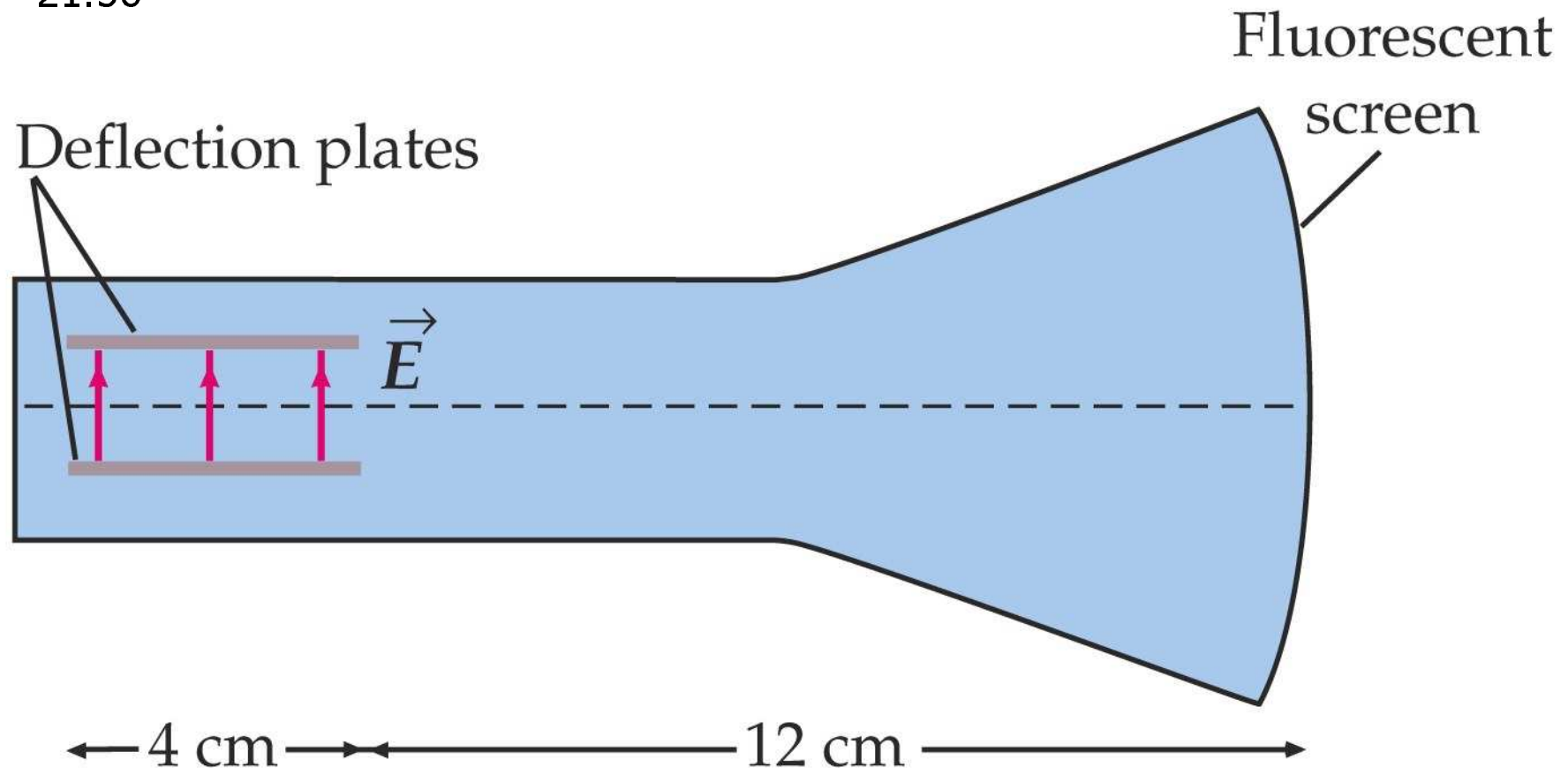
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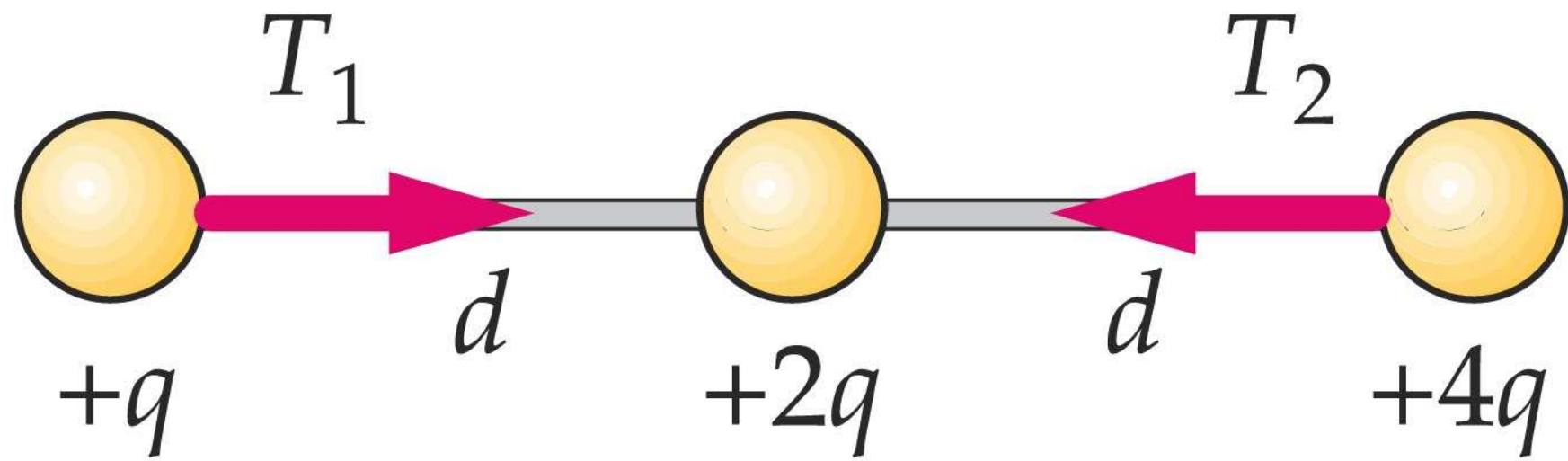
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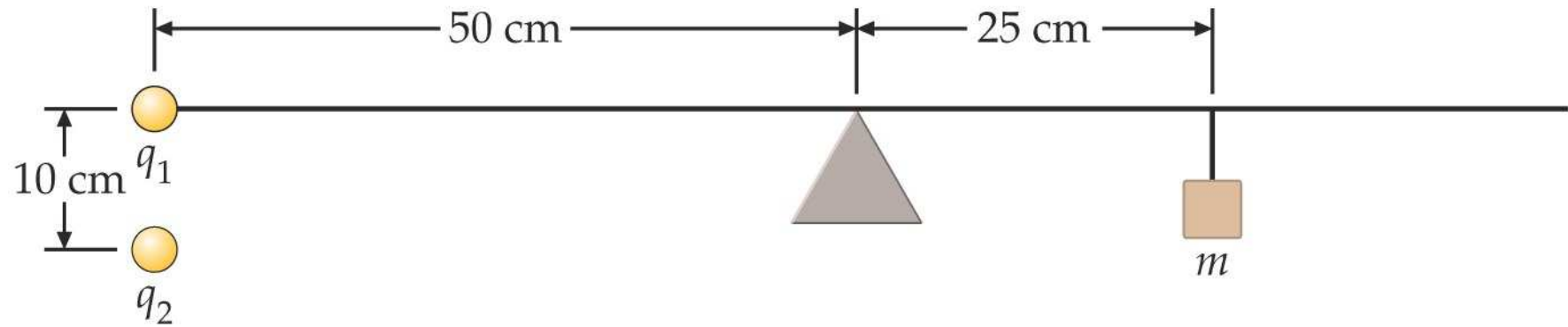
21.56

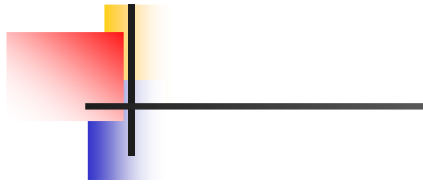


21.68

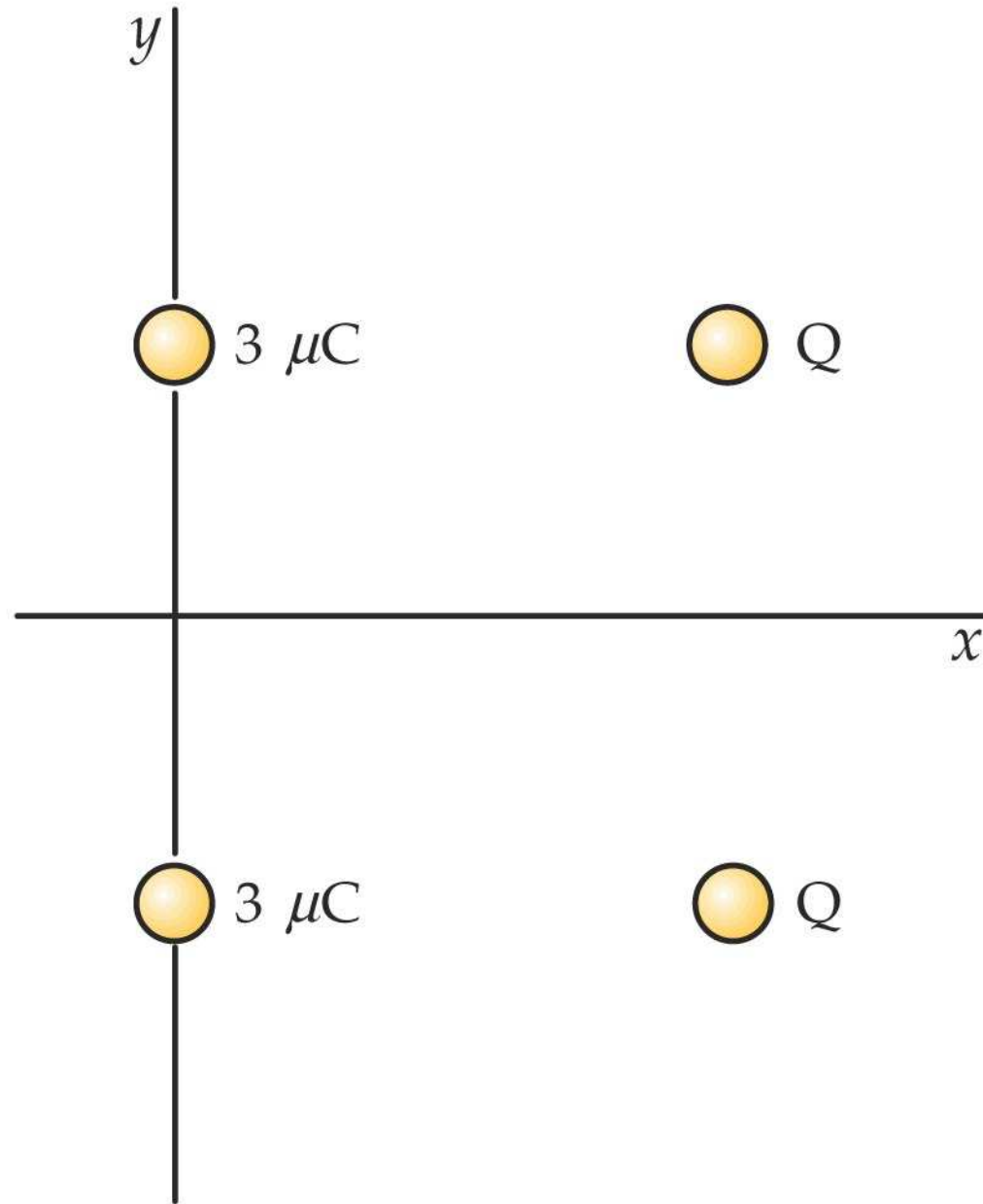


21.73

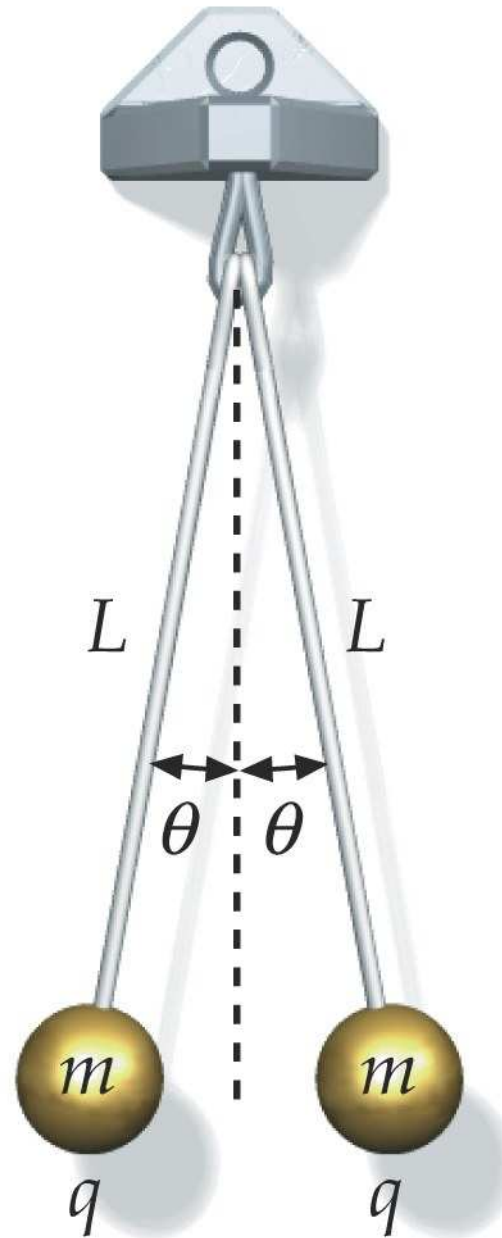




21.74



21.78

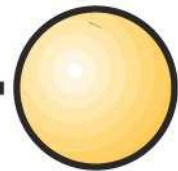




21.80 $-q$



$+q$

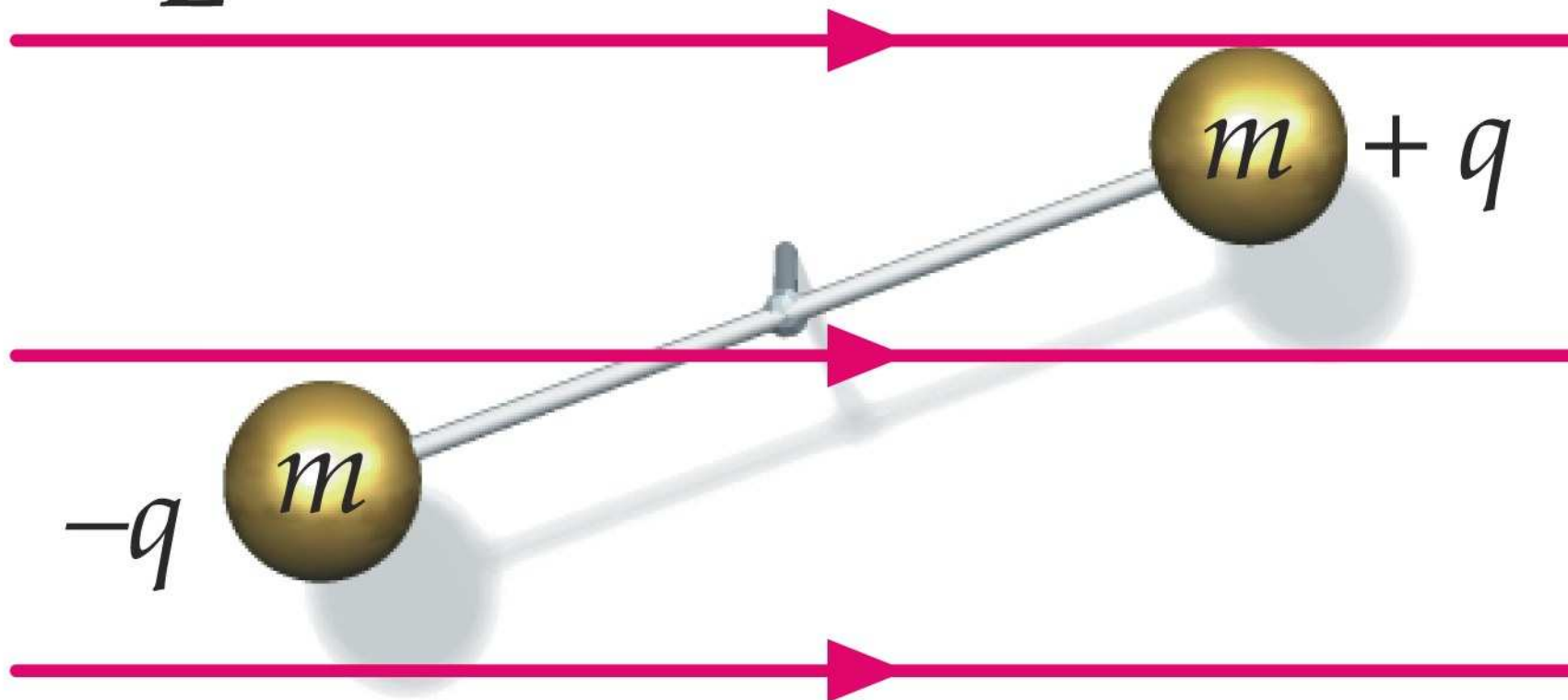


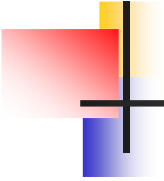
$-q$

$+q$



21.81 $\vec{E}$





21.85



